



Avalanches in fatigue and crossover from creep to fatigue

T. Mäkinen^{1,2}, I. V. Lomakin¹, J. Lahikainen¹, S. Coffeng¹,
M. Tuokkola¹, K. Widell³, J. Savolainen¹, J. Koivisto¹,
and M. J. Alava^{1,2}

¹Department of Applied Physics, Aalto University, Finland

²NOMATEN Centre of Excellence, National Centre for Nuclear Research, Poland

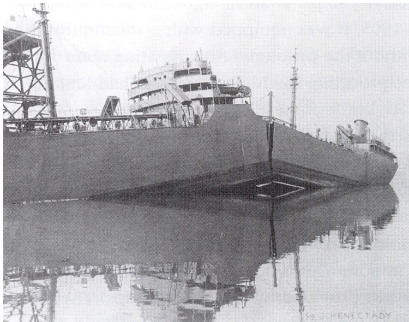
³Department of Mechanical Engineering, Aalto University, Finland

tero.j.makinen@aalto.fi

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What is fatigue?

- | Material response to repetitive loading
- | Important for many practical applications



SS Schenectady

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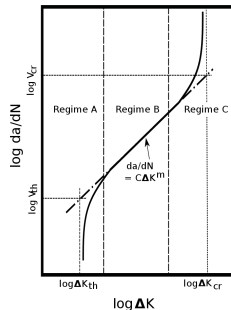


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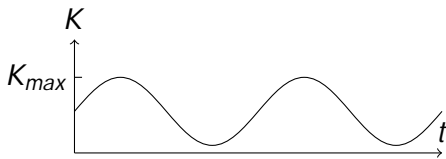
- | Apparent self-similarity (Paris–Erdogan law)

$$\frac{da}{dN} \propto K^m$$

- | a crack length
- | N loading cycles
- | K stress intensity factor
- | $K = K_{max} - K_{min}$



Fatigue protocol

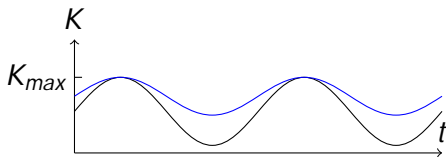


$$R = 0:1$$

| Asymmetry coefficient

$$R = \frac{K_{min}}{K_{max}}$$

Fatigue protocol

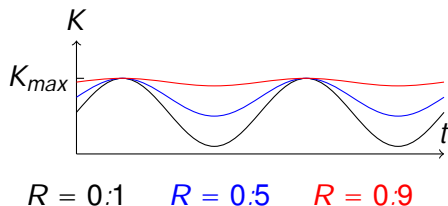


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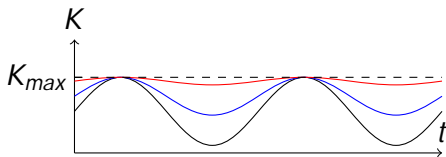
Fatigue protocol



| Asymmetry coefficient

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Fatigue protocol



$$R = 0:1 \quad R = 0:5 \quad R = 0:9$$

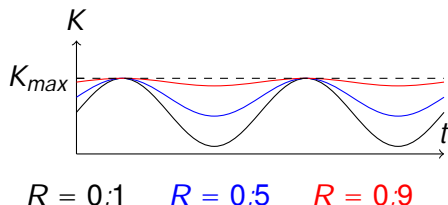
| Asymmetry coefficient

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| Creep as the limiting case

$$R = 1$$

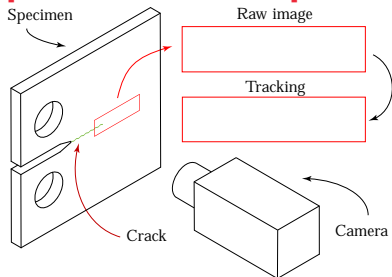
Fatigue protocol



- | Asymmetry coefficient
$$R = \frac{K_{min}}{K_{max}}$$
- | Creep as the limiting case
 $R = 1$

- | the goal here is to explore the intermittent crack growth in fatigue [Kokkonen et al, JSTAT 2017] and the crossover to creep
- | sort of similar to the Oslo Plexiglas experiments [Måløy et al, PRL 2006] but for a geometry more commonly used in the field of material science

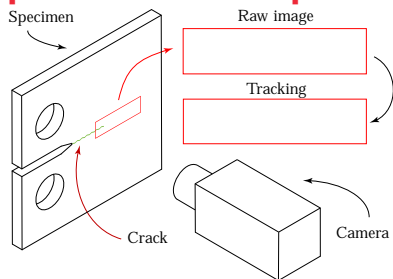
Experimental setup



Lomakin et al, PRRResearch 2021

- | CT-specimens of 1050 Al alloy
- | 10 Hz loading frequency, 0.25 Hz imaging frequency
- | crack tip position tracked from the images

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- | placing the camera at an angle allows the tracking of a fracture line
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Naive attempt from creep to fatigue

Thermally activated process $v = v_c \exp(-E/kT)$

Below depinning threshold $E = \frac{1}{d} \frac{G}{G_c} K^2$

Integrate the velocity over the loading cycle

$$v_{\text{fat}} = \frac{a}{N} = \int_0^{2\pi} v dt = v_c \int_0^{2\pi} \exp\left(-\frac{K(t)}{K_c}\right) dt$$

where $K(t) = \frac{K_{\text{max}} + K_{\text{min}}}{2} + \frac{K}{2} \sin(2\pi t)$

Saddle-point integration gives $v_{\text{fat}} \sim v_c \exp\left(-\frac{K_{\text{max}}}{K_c}\right)$

Aluminum samples

- | flattening of the Paris regime in the Paris plots
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- | power-law distributed jump sizes

$$p(a) \propto a^{-2} e^{-\frac{a}{a_0}}$$

PMMA samples – local velocities

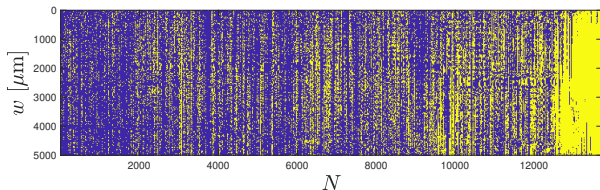
- | velocity of a 1D crack line
- | a power-law distributed tail with an exponent of roughly 2
- | not much variation in the distribution

PMMA samples – local velocities

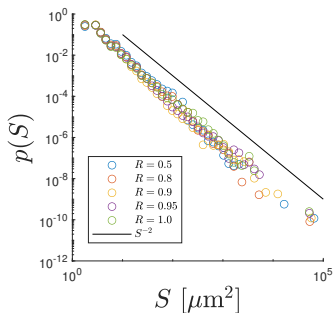
- | velocity of a 1D crack line

- | a power-law distributed tail with an exponent of roughly 2
- | not much variation in the distribution
- | Oslo experiments [Måløy et al, PRL 2006] had an exponent 2.55 (for the normalized velocity)

PMMA samples – avalanche sizes



| the local velocity is thresholded to yield avalanche clusters



| still a power-law with an exponent of roughly 2

| again not much variation in the distribution

| Oslo experiments [Måløy et al, PRL 2006] had an exponent around 1.5-1.7 (some variation depending on the methodology)

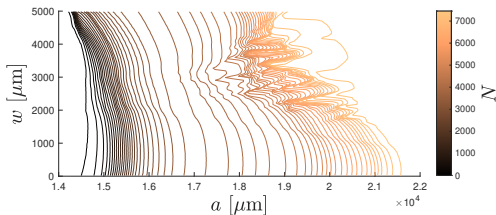
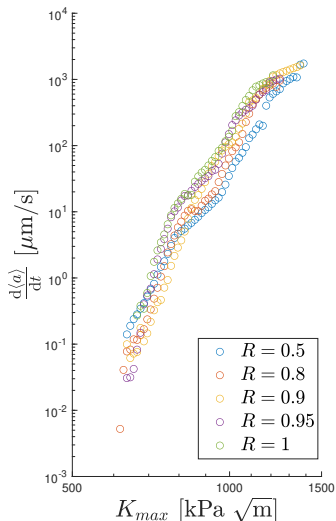
PMMA samples – Paris curves

- | the crack velocity considered as the time derivative of the mean crack position $h a_i$
- | there is a lot of sample-to-sample variation in this mean velocity
 - | the variation seems to increase with increasing R
- | two regimes
 - | related to the strain rate sensitivity of PMMA?

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- | there is a lot of sample-to-sample variation in this mean velocity
 - | the variation seems to increase with increasing R
- | two regimes
 - | related to the strain rate sensitivity of PMMA?
 - | the naive brittle approach works for the end part of the experiments

PMMA samples – front shape



- | the large velocity fluctuations might be related to variations in the front shape
- | pinning at some defect

Conclusions

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- | our plan is to repeat same studies with glass

- | somehow taking the front shape into account might also be good

- | the local dynamics seem fairly constant with varying loading conditions

- | the "universal" exponent of 2 is high
 - | this might partially be explained by the nonstationarity of the velocity