

# Scaling of average avalanche shapes for acoustic emission during jerky twin boundary motion in single-crystalline $\text{Ni}_2\text{MnGa}$

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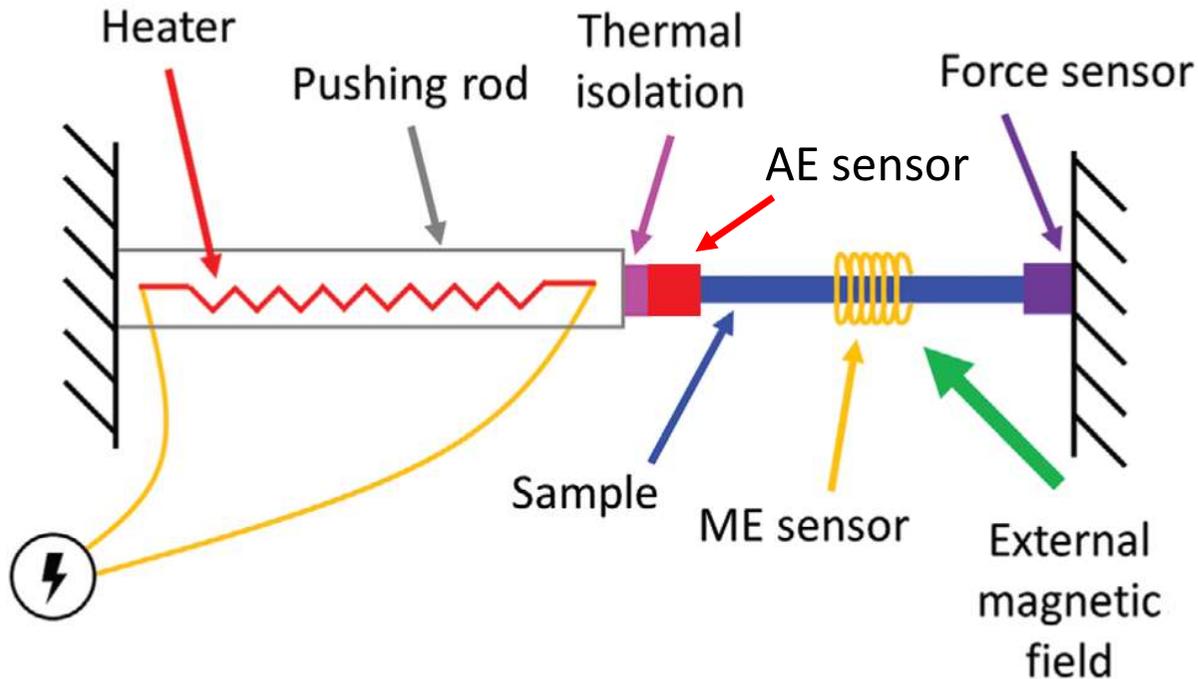
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# Experimental details



## Overview

- 10M modulated  $\text{Ni}_{50.0}\text{Mn}_{28.5}\text{Ga}_{21.5}$  single crystal  
Dimensions 20 mm × 3 mm × 2.5 mm
- Room temperature: fully martensitic state
- Slow compression: 2  $\mu\text{m}/\text{sec}$  rate, using thermal expansion of an alumina rod
- External magnetic field: 0.08 T

## Noise data acquisition:

- Magnetic coil (ME) + Piezoelectric sensor (AE)
- National instruments PCI6111 DAQ board
- 4 MSample/sec

## Stress data acquisition:

- Kistler 9215A (force sensor) + 5015 (amplifier)
- National instruments PCI6071E DAQ board
- 10 kSample/sec

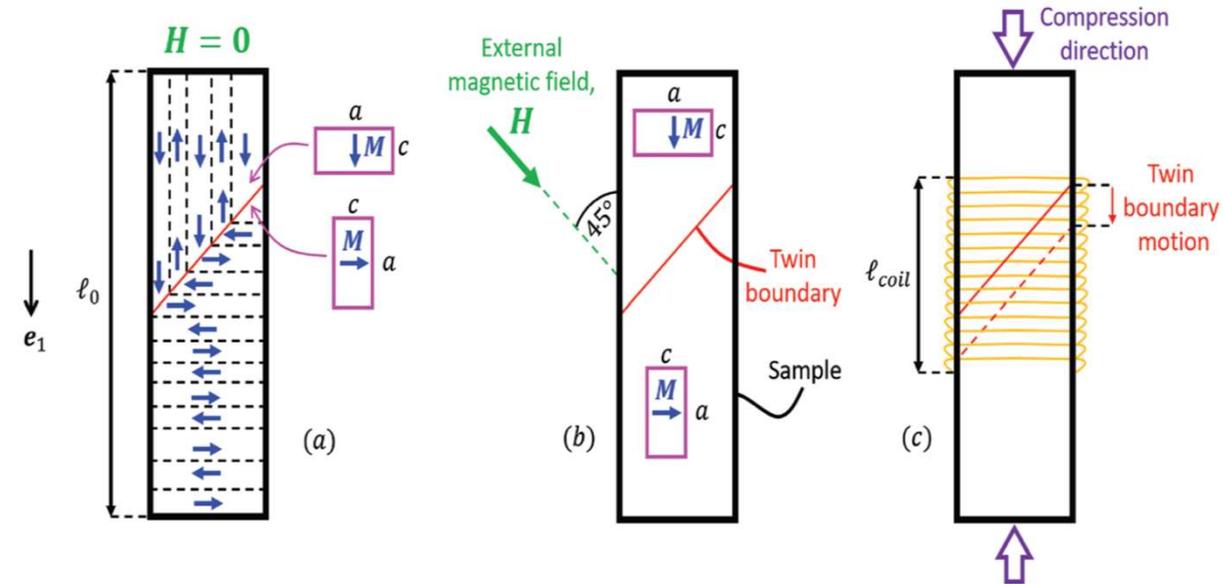


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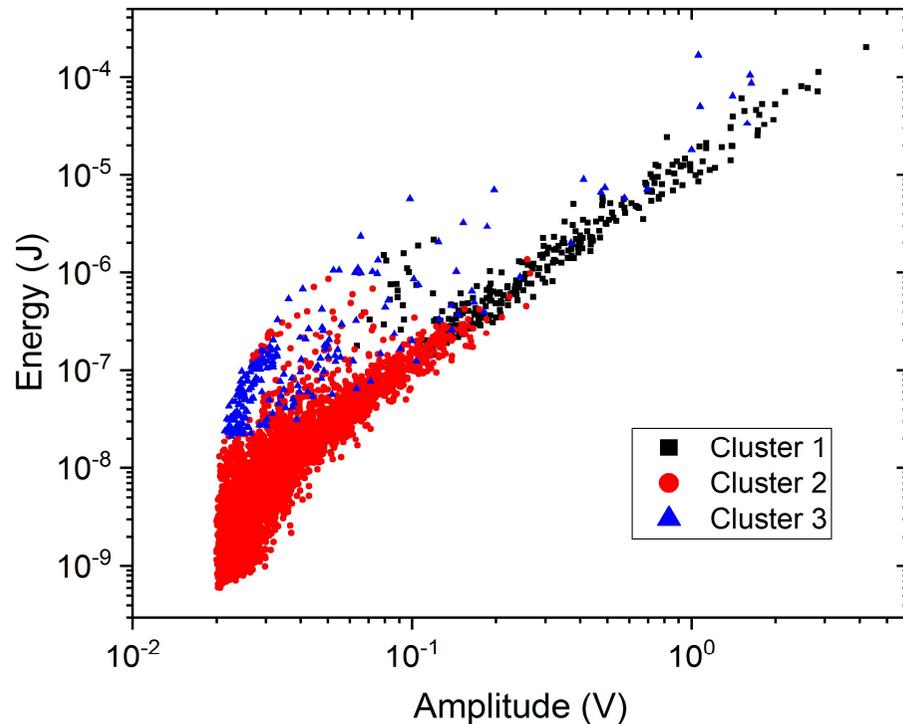
# Experimental details



[1]: E. Bronstein, L. Z. Tóth, L. Daróczi, D. L. Beke, R. Talmon, & D. Shilo: Tracking Twin Boundary Jerky Motion at Nanometer and Microsecond Scales. *Advanced Functional Materials*, **31**(50), 2106573 (2021)

- Investigation of **single twin boundary motion**, induced by **slow compression**; the **magnetizations of the two neighboring twins ensure no magnetic domains at the twin boundary**
- Simultaneous measurement of **stress**, **acoustic emission (AE)** and **magnetic emission (ME)** to study the connection between them.
- There is a **definite agreement between the transformed volumes calculated from the force and the ME measurements** → Ref. [1]
- **This presentation focuses on the acoustic emission results and on the average avalanche profiles and on the scaling relations**

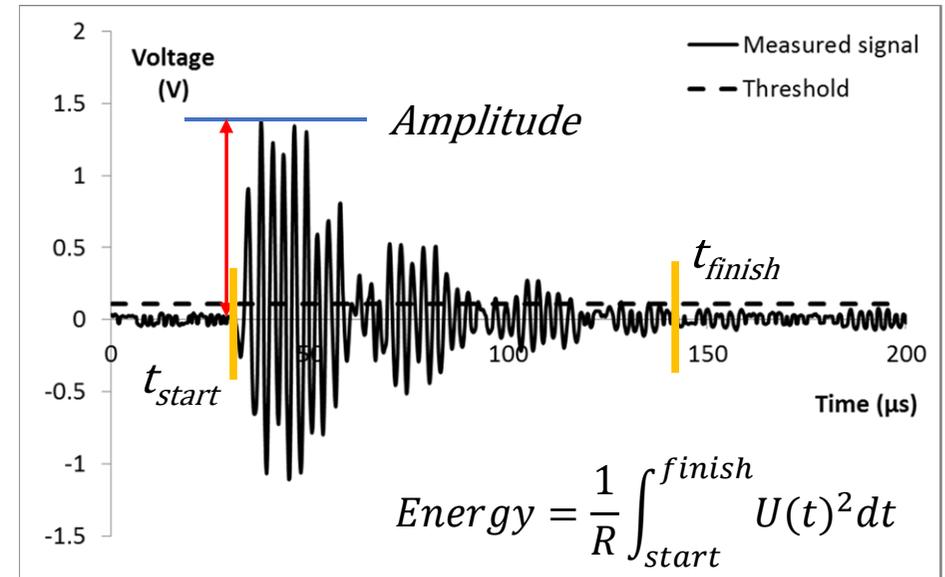
# Processing the AE signals



[2]: E. Pomponi, & A. Vinogradov, *Mechanical Systems and Signal Processing*, **40(2)**, 791-804. (2013)



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- Identification of AE avalanches
- Calculation of AE avalanche parameters
- Clustering [2] to eliminate false AE avalanches and background noise  $\rightarrow$  "Cluster 3"
- Statistical analysis

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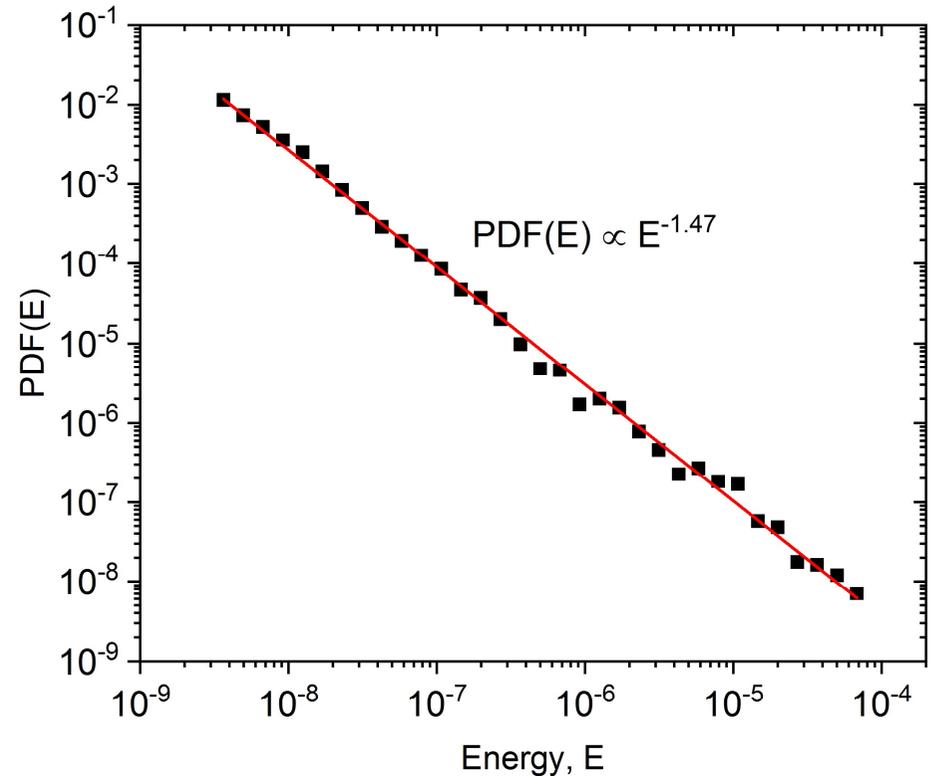
# Acoustic emission results

- Crackling noise (avalanches) with power-law distribution characteristics:

$$P(x) \sim x^{-\eta} \exp\left(-\frac{x}{x_c}\right)$$

- X can be the peak amplitude,  $A_m$ , size, S, or energy, E;  $\eta$  is the characteristic exponent, and  $X_c$  is the cut-off value

Exponent	Magnetic emission [1]	Acoustic emission
Amplitude, $\alpha$	1.6	1.99
Energy, $\varepsilon$	1.26	1.47
Size, $\sigma$	0.8	1.79



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## Average shape of avalanches

- Power-law scaling relations between the above parameters (and between the exponents) were obtained:
- $E \propto A_m^x$ ,  $S \propto A_m^\chi$ , or  $A_m \propto T^\xi$ ;
- In MFT:  $x = 3$ ,  $\chi = 2$  and  $\xi = 1$ .
- The focus is increasingly on the **temporal shapes of avalanches**, i.e. the average of the detected voltage signal ( $U(t)$ ) for a given size (or duration) range of avalanches.
- Average avalanche shapes can be used when the measured voltage is **proportional to the corresponding interface velocity**  $v(t)$ , characteristic for the crackling noise emission.



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## Average avalanche shape for fixed size

The theoretically predicted temporal shape of the avalanche is given by [3]:

$$U(t) = at \exp^{-bt^2},$$

where  $b = 1/\tau^2$ , and  $\tau$  is the characteristic time decay of the avalanche.

The above equation has maxima at:

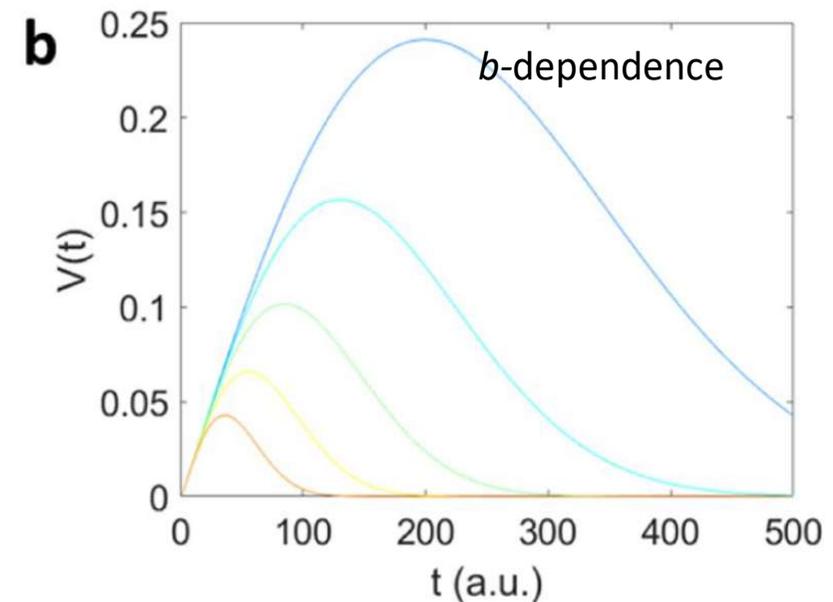
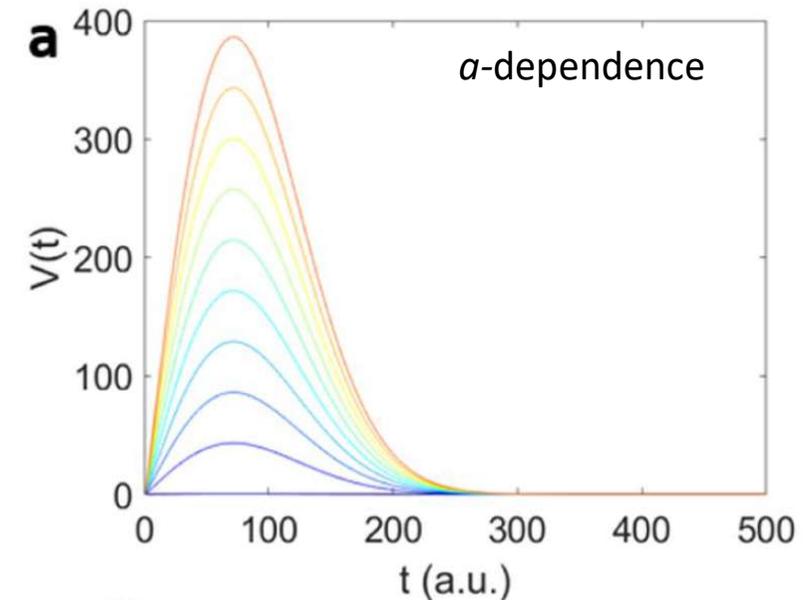
$$t_m = \frac{\tau}{\sqrt{2}} = \frac{1}{\sqrt{2b}}$$

$$A_m = at_m e^{-1/2} = Bt_m$$

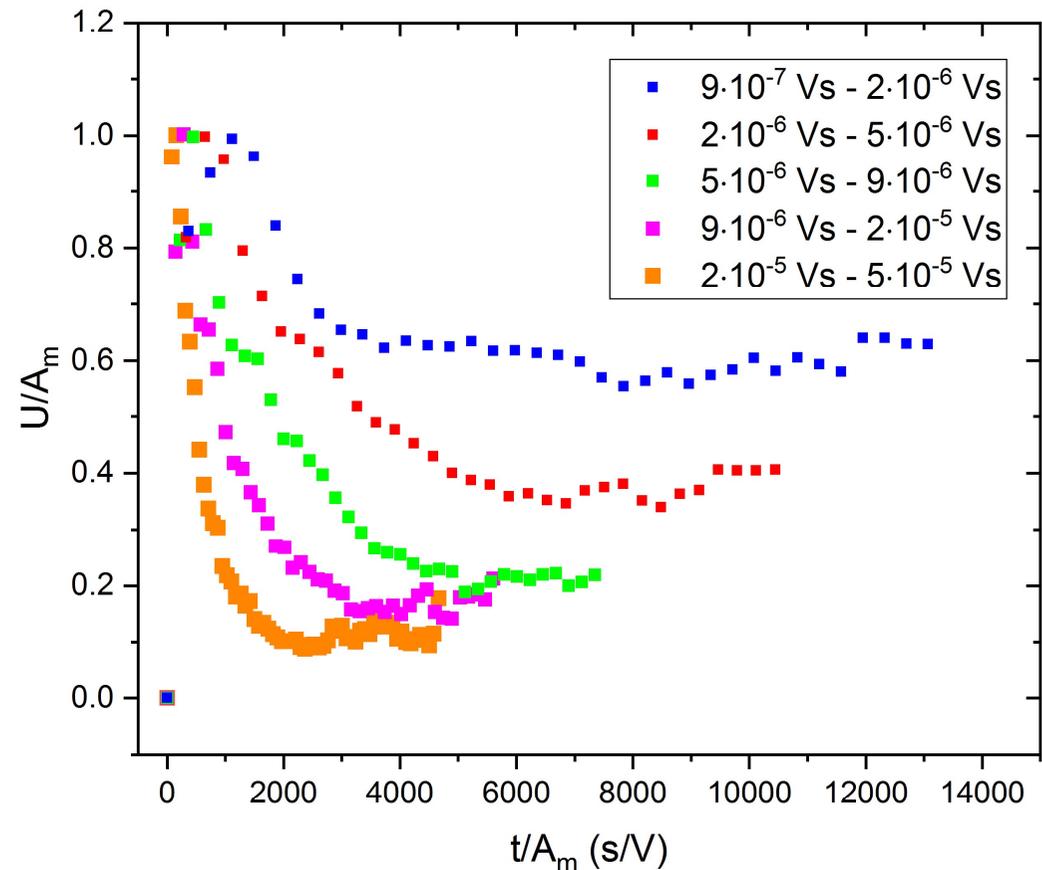


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[3]: Casals, B., *et al. Sci. Rep.* **11**, 5590 (2021).



- **The average avalanche shape** has self-similar behaviour, with an appropriate normalization, they **should be scaled together**, according to a universal scaling function.
- The normalization should be performed with  $A_m$  and  $t_m$  (not distorted by the transfer function of the system [3]).
- In the experiments,  $t_m$  has high uncertainty
- Thus, **if B is constant, we can use  $t_m \propto A_m$**
- **Both axis can be normalized by  $A_m$**



The normalization by  $A_m$  was not working, the scaling was not perfect  
Which implies that **B is not constant**, and thus, a (and b or  $\tau$ ) are not universal constants,  
rather **B has a definite dependence on  $A_m$** [4]:

$$B \propto a \propto A_m^\varphi$$

Thus, using  $A_m = Bt_m$ , we can write:

$$t_m \propto A_m/B \propto A_m^{1-\varphi}$$

Scaling the experimentally determined avalanche shapes with  $A_m$  (voltage axis) and  $A_m^{1-\varphi}$  (time axis) should produce better scaling:

$$U^*(t^*) \propto e^{1/2t^*} \exp - \left( \frac{t^*}{\tau^*} \right)^2 ,$$

where  $U^* = \frac{U}{A_m}$ ,  $t^* = \frac{t}{t_m} = \frac{t}{A_m^{1-\varphi}}$ , as well as  $\tau^* = \frac{\tau}{A_m^{1-\varphi}}$

[4]: S. M. Kamel et al. Materials, 2022,15, 4556



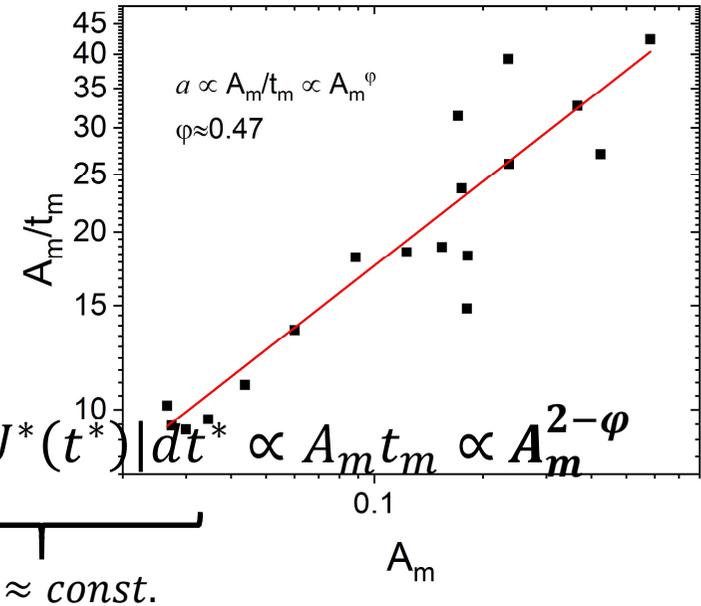
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## How to determine $\phi$ ?

- Directly, from the  $A_m/t_m \propto A_m^\phi$  scaling relation,
- $S \propto A_m^{2-\phi}$ ,
- $E \propto A_m^{3-\phi}$



$$S = \int_0^T |U(t)| dt = \int_0^{T^*} A_m |U^*(t^*)| t_m dt^* = A_m t_m \underbrace{\int_0^{T^*} |U^*(t^*)| dt^*}_{S^* \approx const.} \propto A_m t_m \propto A_m^{2-\phi}$$

$$S^* = \frac{S}{A_m t_m} = \int_0^{T^*} |U^*(t^*)| dt^* = 1.65 \left(1 - e^{-\frac{T^{*2}}{2}}\right) = 1.65 \left(1 - \frac{C}{A_m}\right) \quad T^* \approx \sqrt{2 \ln \frac{A_m}{C}}$$

$T^*$  goes to an asymptotic limit as C (threshold) goes to zero [4].



[4]: S. M. Kamel et al. Materials, 2022,15, 4556

$$\left( U^* = \frac{U}{A_m}, \quad t^* = \frac{t}{t_m}, \quad T^* = \frac{T}{t_m} \right)$$

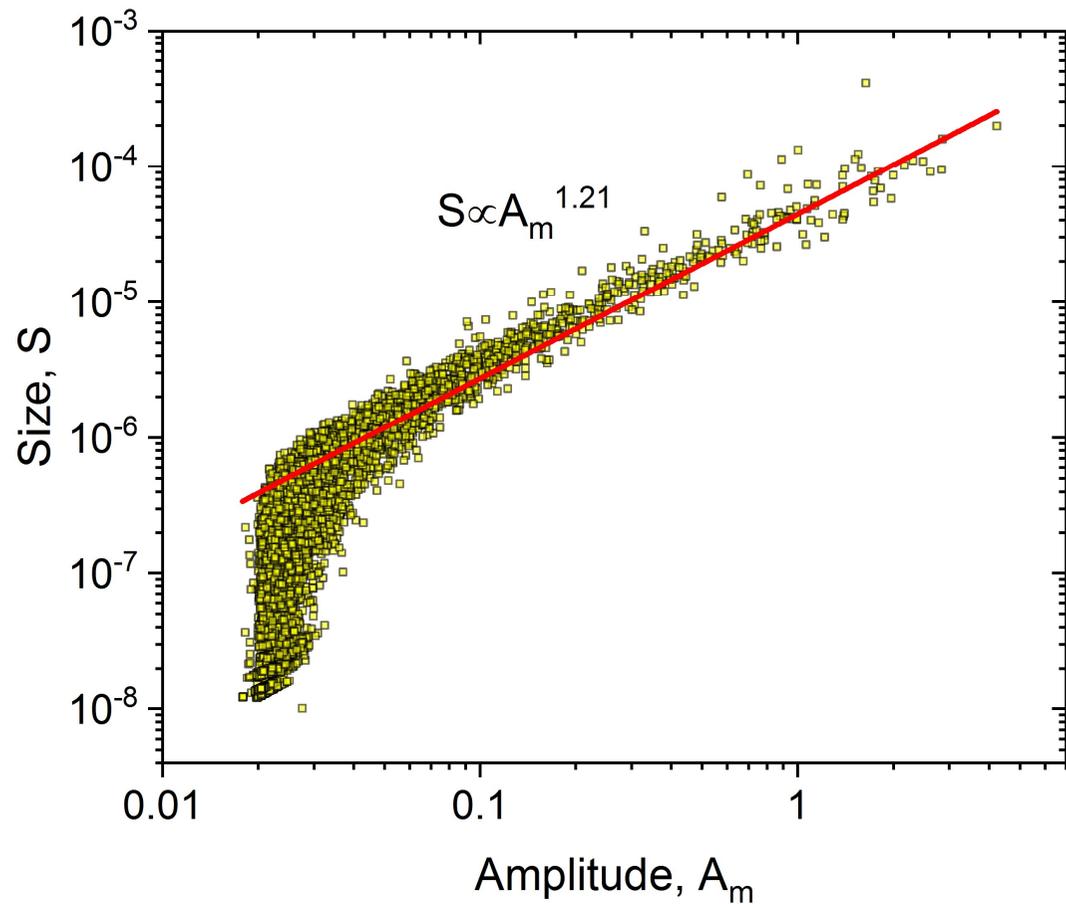
$$S \propto A_m^{2-\varphi}$$

$$2 - \varphi = 1.21$$

$$\varphi_S = 0.79$$

From power exponents of size ( $\sigma=1.79$ )  
and amplitude ( $\alpha=1.99$ ):

$$\frac{\alpha - 1}{\sigma - 1} = 2 - \varphi = 1.24$$

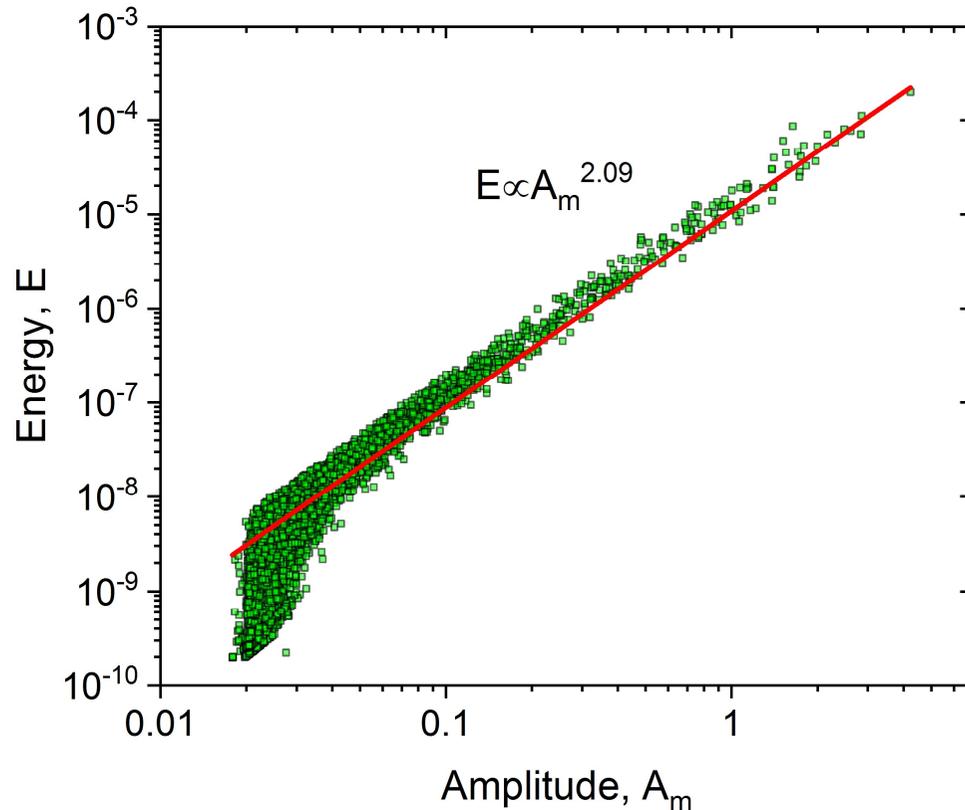


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$$E = \frac{1}{R} \int_0^T U^2(t) dt \propto \int_0^{T^*} A_m^2 U^{*2}(t^*) t_m dt^* = A_m^2 t_m \int_0^{T^*} U^{*2}(t^*) dt^* \propto A_m^2 t_m \propto A_m^{3-\varphi}$$



$$3 - \varphi = 2.09$$

$$\varphi_E = 0.91$$

From power exponents of energy ( $\varepsilon=1.47$ ) and amplitude ( $\alpha=1.99$ ):

$$\frac{\alpha - 1}{\varepsilon - 1} = 3 - \varphi = 2.10$$



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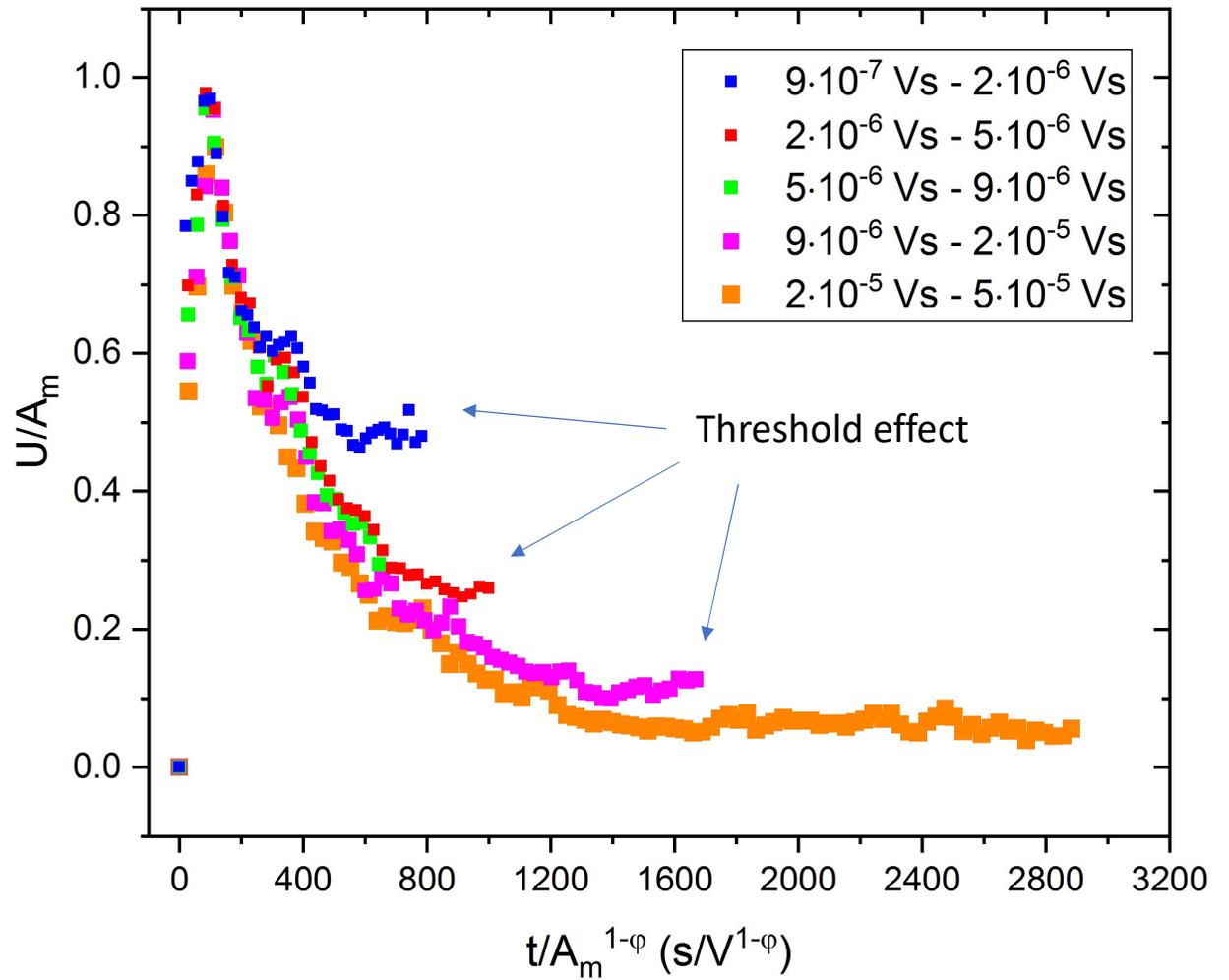
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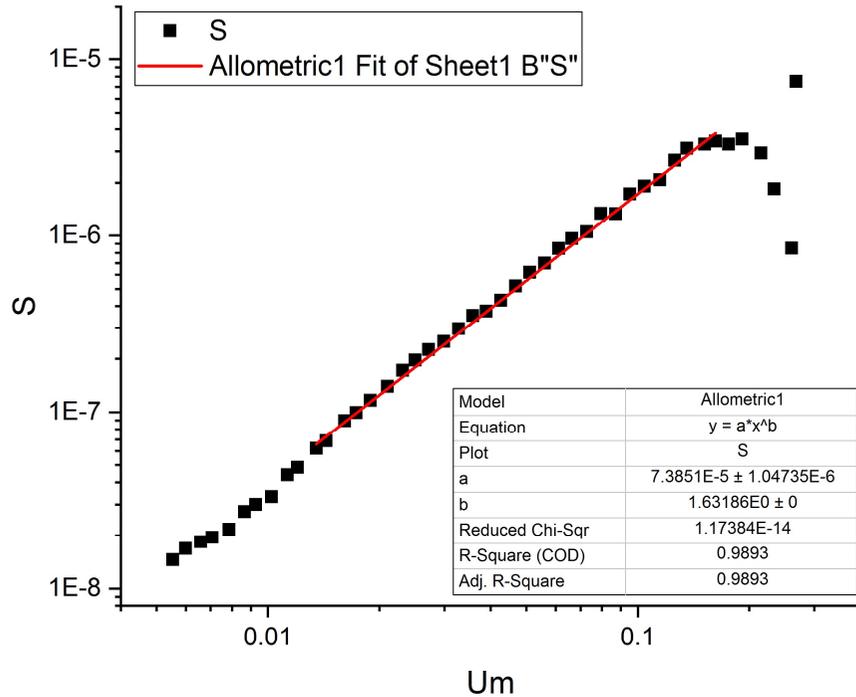
Average avalanche shapes for various fixed avalanche sizes,  
 Normalized by  $A_m$  and  $A_m^{1-\varphi}$ ,  $\varphi=0.85$

$$\varphi = \frac{\varphi_S + \varphi_E}{2} = 0.85$$

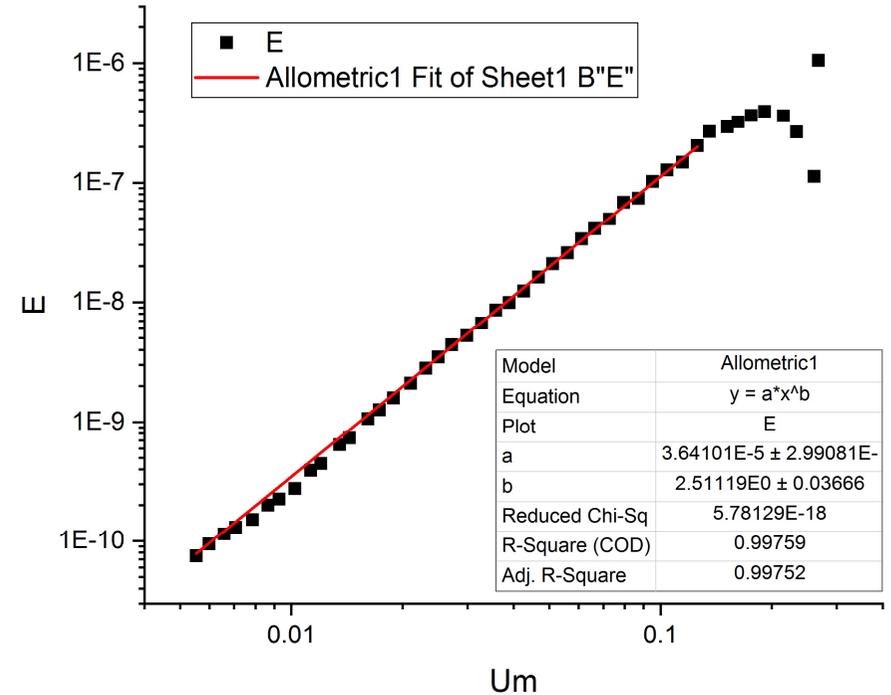
- Impressive scaling of the avalanches belonging to different size ranges.
- The effect of the threshold is visible near the end of the curves (especially for the smallest size range).



# Magnetic emission



$$\frac{\partial \ln S}{\partial \ln U_m} \approx 2 - \varphi \approx 1.63 \quad \varphi_S \approx 0.37$$



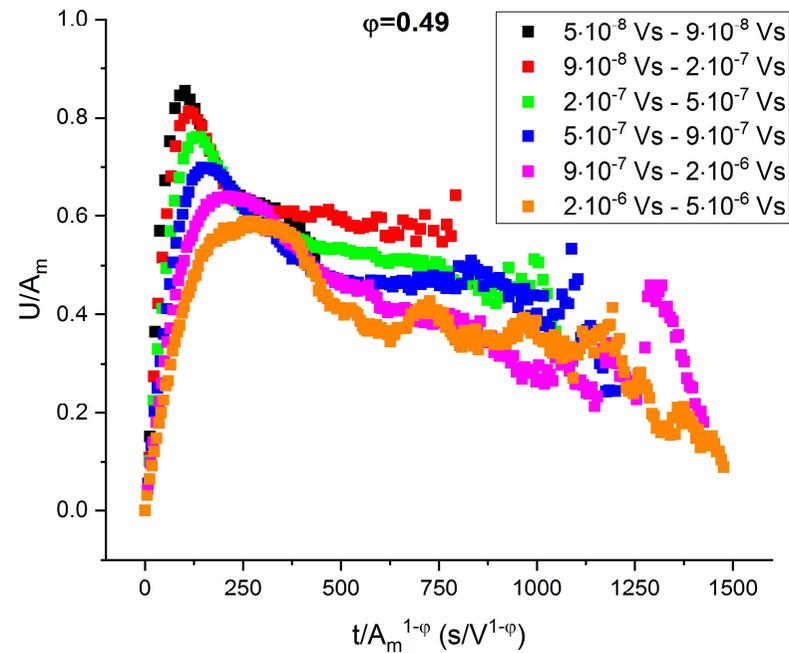
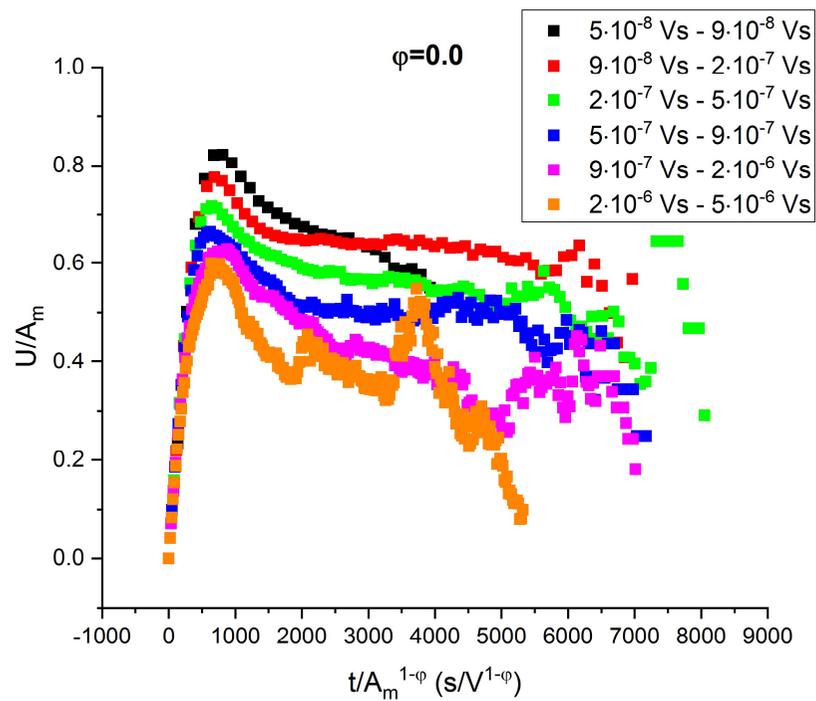
$$\frac{\partial \ln E}{\partial \ln U_m} \sim 3 - \varphi \quad \varphi_E \approx 0.49$$



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- The scaling of the magnetic avalanches is poor
- Experimental reasons: Eddy currents, ...



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## Average avalanche shape for fixed duration [5]:

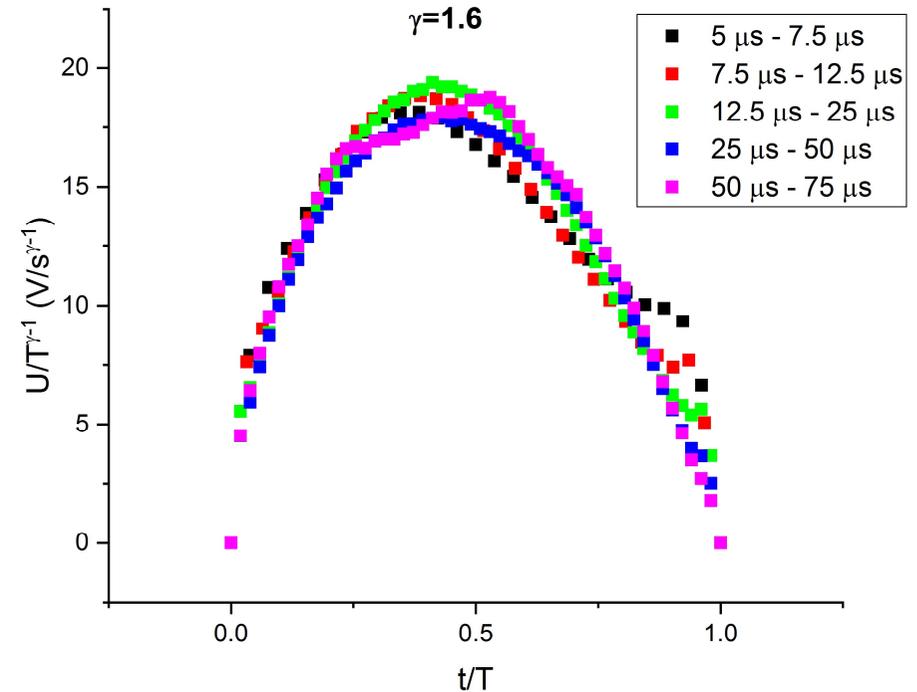
$$\langle U(t|T) \rangle \propto T^{\gamma-1} \left[ \frac{t}{T} \left( 1 - \frac{t}{T} \right) \right]^{\gamma} \left[ 1 - a \left( \frac{t}{T} - \frac{1}{2} \right) \right]$$

$$S \propto T^{\gamma} \quad \longrightarrow \quad \gamma = 1.6$$

Very good scaling of the average shape for various fixed duration ranges

Visible skewness: Eddy currents

[5]: Laurson, et al. "Evolution of the average avalanche shape with the universality class." *Nature communications* 4.1 (2013): 1-6.



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