Géza Ódor EK-MFA Complex Systems Department, Budapest





Energiatudományi Kutatóközpont

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Theoretical research and experiments suggest that the brain operates at or near a **critical state** between sustained activity and an inactive phase, exhibiting optimal computational properties (see: *Beggs & Plenz J. Neurosci. 2003; Chialvo Nat. Phys. 2010; Haimovici et al. PRL 2013*)



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Quasistatic inhomogneity causes dynamical criticality in Griffiths phases



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Quasistatic inhomogneity causes dynamical criticality in Griffiths phases

→ Mixed order transition + Griffiths phase together





Electrode LFP experiments Since *Beggs* and *Plenz 2003* For humans and animals



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Diverging fluctuations → High sensitivity to stimuli



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Cons: Tuning to critical point is needed Danger of super-critical (epileptic) behavior Self-organization to criticality (SOC) ?





Higher order interactions!



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Order parameter : density of active sites (ρ)



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Mean field for reaction diffusion systems : $mA \rightarrow (m+k)A$, $nA \rightarrow (n-l)A$ For m > n : first order phase transition see: GÓ: RMP 76 (2004) 663.



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In case of quenched heterogeneity: Griffiths Phase

PRL 105, 128701 (2010)



• Fixed (quenched) disorder/impurity

changes the local birth rate $\Rightarrow \lambda_{c} > \lambda_{c}^{0}$











contribute to the density: $\rho(t) \sim \int dL_R L_R w(L_R) \exp[-t/\tau(L_R)]$



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saddle point analysis: stretched exponential **Griffiths Phase**

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Act.

Abs.



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- For $\lambda_c^{\ o} < \lambda < \lambda_c$: $\tau (L_R) \sim \exp(b L_R)$: Griffiths Phase $\rho(t) \sim t^{-c/b}$ continuously changing exponents





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- At λ_c : b may diverge $\rightarrow \rho(t) \sim \ln(t)^{-\alpha}$ Infinite randomness fixed point scaling
Rare Region theory for quench disordered CP CP: infect with prob λ , heal with prob 1- λ



 λ_c^{o} "clean critical point" Abs.

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- At λ_{ρ}^{0} the characteristic time scales as: $\tau (L_{\rho}) \sim L_{\rho}^{Z} \Rightarrow$ saddle point analysis: In $\rho(t) \sim t^{d/(d+Z)}$ stretched exponential
- For $\lambda_c^0 < \lambda < \lambda_c$: τ ($L_{\rm p}$) ~ exp(b $L_{\rm p}$): Griffiths Phase $\rho(t) \sim t^{-c/b}$ continuously changing exponents
- At λ_{c} : b may diverge $\rightarrow \rho(t) \sim \ln(t)^{-\alpha}$ Infinite randomness fixed point scaling
- **GP:** Dynamical (scaling) criticality + susceptibility diverges



Diffusion and structural MRI images with 1 mm³ voxel resolution : 10 ⁵-10 ⁶ nodes



Diffusion and structural MRI images with 1 mm^3 voxel resolution : $10^5 - 10^6$ nodes

Hierarchical modular graphs



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fiber \rightarrow edge (~ 10¹⁰)

+ noise reduction \rightarrow graph undirected, weighted









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+ noise reduction → graph
undirected, weighted
Graph dimension: d < 4 *MG*, *GO Sci.Rep.* 2016









KKI-18 graph: 836733 vertex, 8 x 10⁷ weighted, undirected edges



FIG. 1. Link weight probability density function of the KKI-18 OCP graph. Dashed line: a PL fit for intermediate w_{ij} 's. Inset: survival probability in the K = 6 threshold model near the transition point for $\lambda = 0.003$, $\nu = 0.3, 0.4, 0.45, 0.5, 0.55, 0.6, 0.7$ (top to bottom curves).

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FIG. 2. Avalanche survival distribution of the relative threshold model with K = 0.25, for $\nu = 0.95$ and $\lambda = 0.8, 0.81$, 0.82, 0.83, 0.835, 0.84, 0.845, 0.85, 0.86, 0.87, 0.9, 0.95, 1 (bottom to top curves). Inset: Local slopes of the same from $\lambda = 0.835$ to $\lambda = 1$ (top to bottom curves). Griffiths effect manifests by slopes reaching a constant value as $1/t \rightarrow 0$.



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Autocorrelations show the same \rightarrow "Burstyness"



 $p(s) \propto s^{-\tau}$



FIG. 3. Avalanche size distribution of the relative threshold model with K = 0.25, for v = 1 and $\lambda = 1,0.9,0.8$. Dashed line: PL fit to the $\lambda = 0.8$ case. Inset: Avalanche shape collapse for T = 25,63,218,404 at $\lambda = 0.86$ and v = 0.95.

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Scaling near experimental values in the Griffiths Phase (*GO PRE 2016***)**

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Addition of a third (refractive) state does not destroy GP



FIG. 3: Avalanche size distribution in the relative threshold model with refractory states, for K = 0.2, $\nu = 1$ and $\lambda = 0.91, 0.965, 0.985, 0.995$ (bottom to top symbols). Lines: PL fits for $10^2 < s < 10^5$, for these curves as shown by the legends.

Robustness of Griffiths effects in homeostatic connectome threshold models *G. Ó, Phys. Rev. E 98 (2018) 042126*

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Time dependent threshold model : GP shrinks, but survives for weak variations





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FIG. 10: Avalanche size distribution of the time dependent relative threshold model with 30% inhibitory links at K = 0.1, $\Delta K = 0.01$, $\nu = 0.95$ and $\lambda = 0.473, 0.478, 0.480, 0.493$ (bottom to top symbols) Dashed lines: PL fits for the tails of the $\lambda = 0.473$ and $\lambda = 0.493$ curves (bottom to top).

HMN2d:

Exponentially decaying connection probabilities with the levels l:

 $p_l \approx \langle k \rangle (1/2) sl$



FIG. 1: Two lowest levels of the hierarchical network construction with 4 nodes/module. Dashed lines: $l = l_1$, dotted lines: $l = l_2$. The solid lines denoted R1 are randomly chosen connections within the bottom level (fully connected) modules, while those denoted R2 provide random connections on the next level. Links can be directed.

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where **Griffiths Phase** is present





->





FIG. 3: Number of nodes within chemical distance r in HMN2d networks with s = 4 and l = 9 levels. Different curves correspond to different $\langle k \rangle$ -s. Inset: local slopes d_{eff} of the N(r) curves, defined in Eq. 4.



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10¹

r

10⁰

10⁰

2.2 2.6 2.8 3 3.03

- 3.06 - 3.1 - 3.2

10²

Topological dimension : $N(r) \sim r d$ **Effective dimension:** $d_{eff} = \frac{\ln[N(r))/N(r')]}{\ln(r/r')}$ **Breadth-first search results,** in agreement with the
1d networks with power-law ranged, long edges:For $s = 4 : \langle k \rangle$ dependent continuously changing
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For *s* < 4 small-world networks, $d \rightarrow \infty$



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We study s = 3 *now* in more detail

+ lattice connectedness at: l = 1



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$$d_{eff} > 4 \rightarrow$$
 mean-field behavior expected !

Two-state system: $x_i = 0, 1$ (inactive, active)

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Mean-field approximation: probability of site activation: ρ and a pair of nodes can be selected in a (*N*-1)(*N*-2)/2 way. The creation rate is: $\frac{1}{2}(N-1)(N-2)\Lambda\rho^2(1-\rho)$

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Calling: $\lambda = (N - 1)(N - 2) \Lambda/2$, for a full graph of *N* nodes

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♀ ,absorbing phase	
1/2 _ active phase	
	λ
$0 \qquad \lambda_c = 4/5$	

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By running the model on the HMN2d graphs : discontinuous transition + PLs with continuously varying exponents



Fig. 1. Evolution of the activity for different control parameter λ in case of start from fully active state in a threshold model running on hierarchical modular graphs. One can see a discontinuous transition to a Griffiths Phase (from Ref. [4]).



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 $\rho(t) \sim \int dn \ n \ p(n) \ \exp\left[-t/\tau(n)\right]$



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For $\lambda_c^{0} < \lambda < \lambda_c : \rho(t) \sim t \cdot c/b$

MEASUREMENTS

- Density of active nodes $\rho(t) = 1/N \sum_{i=1}^N x_i$
- A single pair of active nodes can trigger an avalanche of duration T and spatiotemporal size $s = \sum_{i=1}^{N} \sum_{t=1}^{T} x_i$. It allows us to compute:
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FIG. 3. Avalanche size distributions at different λ branching rates, denoted by the symbols, in the presence of excitatory links in the HMN2d with l = 5, 6 levels. From top to bottom curves: $\lambda = 0.33, 0.325, 0.322, 0.32$ (l = 5 cyan and l = 6 green), 0.315, 0.31. Dashed lines show PL fits for the tails: s > 1000 at $\lambda = 0.315, 0.322, 0.32$.

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FIG. 4. Survival probability of the activity at different branching rates in the K = 2 threshold model with excitatory links. From top to bottom curves: $\lambda = 0.33, 0.325, 0.322, 0.32 \ (l = 5 \text{ and } l = 6), 0.315$. Dashed lines show PL fits for the tails: $s > 10^4$ at $\lambda = 0.315, 0.32$, 0.322, 0.33.

Explanation for the Griffiths Phase

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Hubs or cores in modules remain active which decay as:

 $ho \sim \exp(-t/\tau_{
m ls})$

Random, inter-module connections with single links \leftrightarrow K=2 \rightarrow quasi unconnected, finite rare regions

Conclusions

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- Griffiths phase (GP) can occur in high dimensional systems due to fragmentation of the activity propagation caused by the modules
- Nonuniversal PLs suggest that Griffiths effects are present



Fig. 5: Steady-state behavior for the excitatory, inhibitory, and refractory-inhibitory cases. Inset: evolution of ρ in an inhibitory HMN2d with N = 4096 for different initial activity densities: $\rho(0) = 0.0005, 0.001, 0.01, 0.1, 1$ (bottom to top curves).

Discontinuous jump in ρ, metastability and GP: Hybrid Phase Transition!

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 Discontinuous jump in ρ, metastability and GP: Hybrid Phase Transition! May apply to other heretogeneous, excitable systems G. Ó. & B. S. PHYSICAL REVIEW RESEARCH 3, 013106 (2021)

Thank you for your attention !
Inhibitory links (10-30%) generate Griffiths Phase



FIG. 5. Avalanche survival distribution of the relative threshold model with 30% inhibitory links at K = 0.1, for $\lambda = 0.95$ and $\nu = 0.4, 0.45, 0.49, 0.5, 0.51, 0.52, 0.550, 57, 0.7$ (bottom to top curves). Inset: Local slopes of the same curves in opposite order.

Inhibitory links (10-30%) generate Griffiths Phase with non-universal power laws and ultra-slow dynamics at λ_c

 $\tau \sim 1.3 - 2$



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FIG. 6. Avalanche size distribution of the relative threshold model with 30% inhibitory links at K = 0.1, $\nu = 0.95$, and $\lambda = 0.49, 0.5, 0.55$. Dashed lines: PL fits. Inset: Effective η exponent for $\nu = 0.95$ and $\lambda = 0.49, 0.5, 0.51, 0.51, 0.55$ (bottom to top curves).

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G.Ó. Phys. Rev. E 94, 062411 (2016)

Inhibitory (negative) links compared to experiments

Inhibitions: 20% of links: $w_{ij} \rightarrow -w_{ij}$ randomly

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Duration scaling exponent within experimental range: $1.5 < \tau_t < 2.4$ *J.M. Palva et al PNAS 110 (2013) 3585*

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FIG. 6. Survival probability of the activity at different branching rates and $v = 1 - \lambda$ for the K = 2 threshold model with levels: l = 5, 6 for the case with 20% of inhibitory links. From bottom to top symbols: $\lambda = 0.5, 0.505, 0.510, 0.515, 0.520$ (l = 5 purple cross and l = 6 blue circle), 0.525 (l = 5 brown cross and l = 6 brown circle). Dashed lines are PL fits for the tails of $\lambda = 0.505, 0.525$ curves.

Inhibitory refractory: nodes cannot be reactivated for Δt time

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FIG. 7. Avalanche size distributions at different λ branching rates, denoted by the symbols, in case of the refractory model, in the presence of inhibitory links in HMN2ds with l = 5, 6 levels. From bottom to top symbols: $\lambda = 0.39$, 0.40 (l = 5 left triangle and l = 6up triangle), 0.41, 0.42, 0.43. Dashed lines are PL fits for the tails of $\lambda = 0.39$, 0.4, 0.41, 0.43 cases for t > 1000. The inset shows the oscillatory behavior of $\rho(t)$ of a single run for $\Delta t = 10$.

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FIG. 8. Survival probability of the activity at different branching rates λ for the levels l = 5, 6, in the case of the inhibitory-refractory model. From bottom to top symbols: $\lambda = 0.40$, 0.41 (l = 5 and l = 6), 0.42 (l = 5 light green and l = 6 dark green), 0.43. Dashed lines show PL fits for t > 1000 for the $\lambda = 0.4$, 0.41, 0.43 cases. Inset: $\rho(t)$ at $\lambda = 1$, l = 7 averaged over 10⁵ realizations. Blue boxes: excitatory; red diamonds: inhibitory. Black bullets: BFS $\rho(r)$ results. Dashed lines are PL fits for the initial regions: $1 \le t <$ 10) resulting in effective dimensions: $d_{\text{eff}} = 1.84(3)$ (excitatory), $d_{\text{eff}} = 1.19(1)$ (inhibitory), d = 4.18(5) (graph dimension estimated for 5 < r < 10).